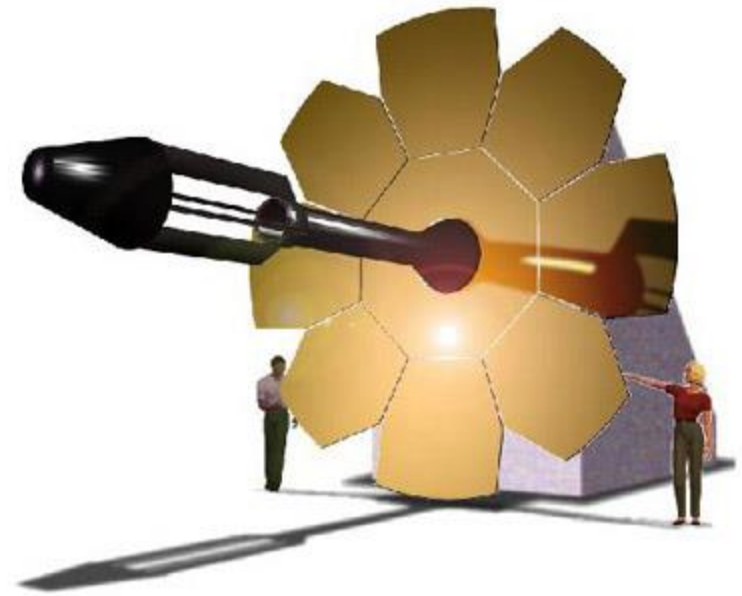




Reflector Telescope Design

Objective: Describe the optical imaging performance of reflector telescopes. Telescope eyepieces are not discussed here, but must be designed carefully to not destroy the imaging capabilities of a telescope.

- + mirrors have no chromatic aberration (same telescope useful in uv-vis-IR)
- + mirrors have high reflectance over a very broad wavelength range
- + large mirrors can be built stronger and lighter than large lenses
- mirrors easily get in the way of each other





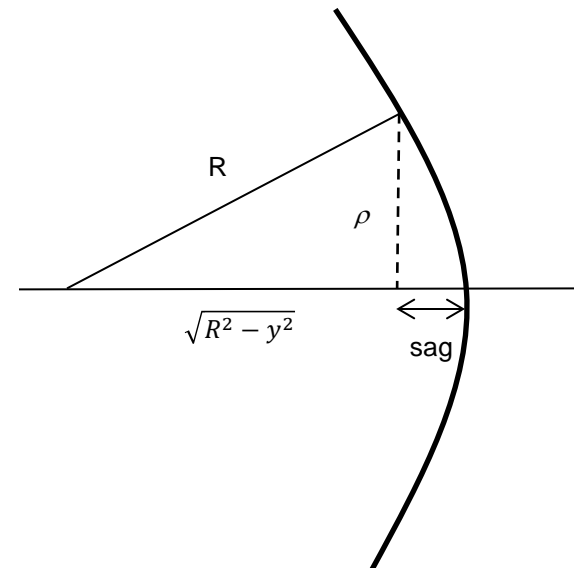
Surface sag

“sag” is the optical term for the shape of a surface that deviates from flat.

$$\text{Sag} = R - \sqrt{R^2 - \rho^2} = R - R \left[1 - \left(\frac{\rho}{R} \right)^2 \right]^2$$

$$\text{Sag} \cong R - R \left[1 - \frac{\rho^2}{2R^2} \right] = \frac{\rho^2}{2R}$$

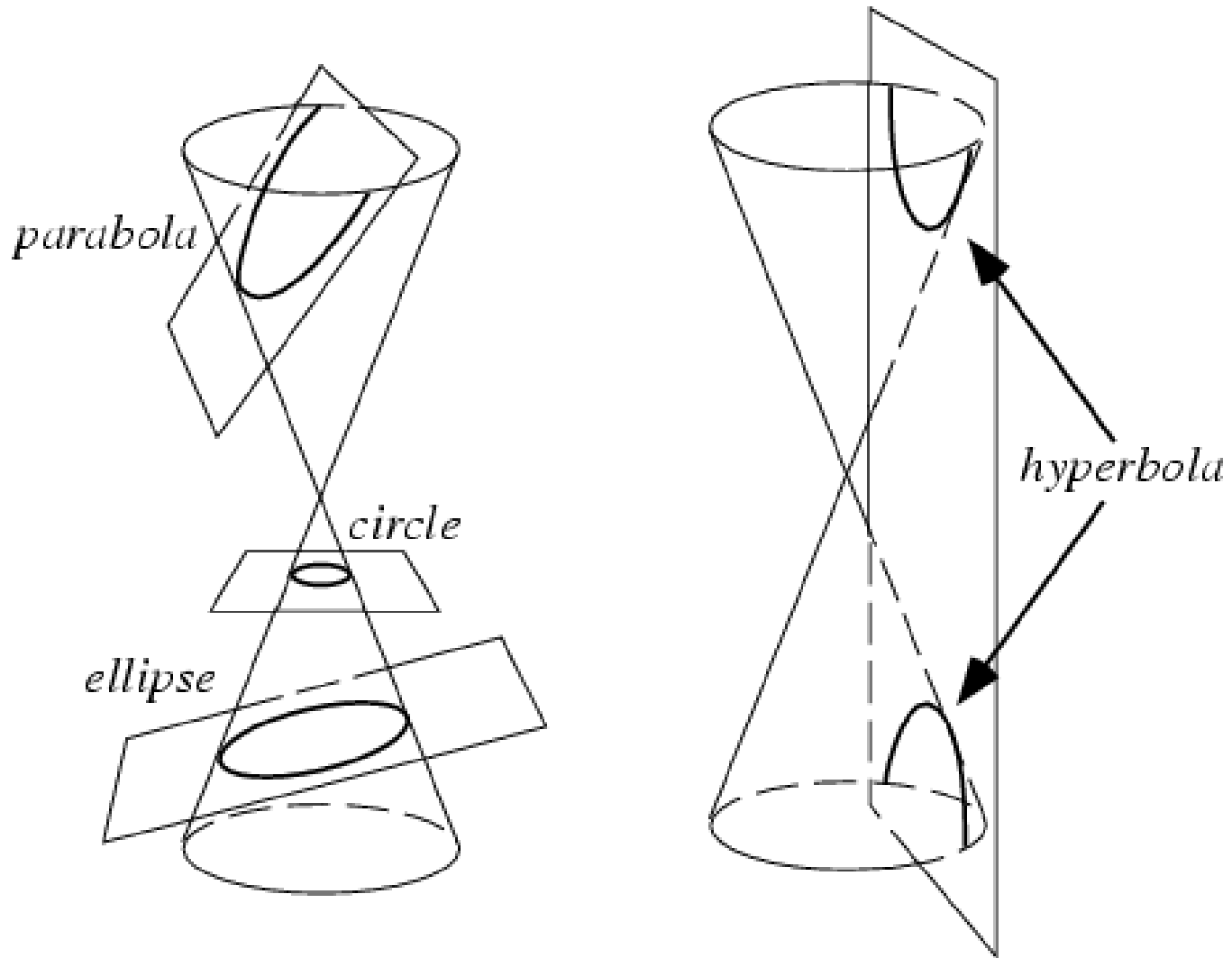
$$\text{Sag} \cong \frac{\rho^2}{2R}$$



...parabolic approximation of the sag of a spherical surface...



Conic Sections





Conics

A conic is a surface of revolution formed by spinning a conic section around the axis.

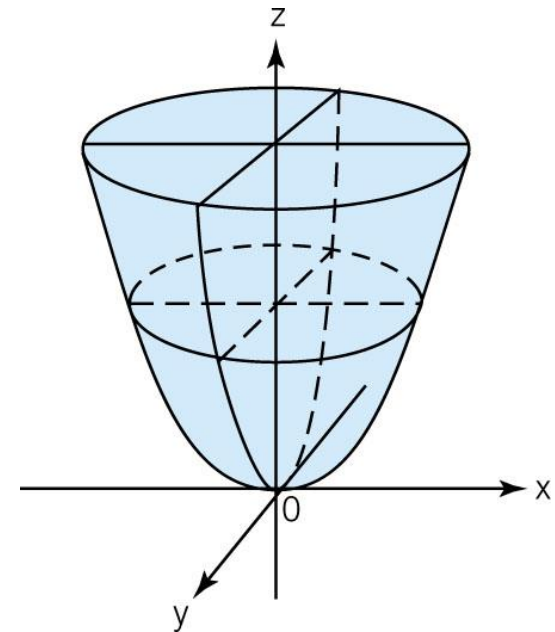
Equation for a conic centered on the z axis: $\rho^2 - 2Rz + (k + 1)z^2 = 0$

ρ = radial coordinate ($\rho^2 = x^2 + y^2$)

R = vertex radius of curvature,

k = conic constant.

All conics satisfy Fermat's principle (that is, have perfect imaging) at 2 conjugate points.





Conic constants

conic constant = K

Eccentricity = e

$$k = -e^2$$

| <u>Shape</u> | <u>eccentricity</u> | <u>conic constant</u> | <u>focal length</u> |
|-------------------|---------------------|-----------------------|---------------------|
| sphere | $e = 0$ | $k = 0$ | $f < R/2$ |
| paraboloid | $e = 1$ | $k = -1$ | $f = R/2$ |
| Hyperboloid | $e > 1$ | $k < -1$ | $f > R/2$ |
| Prolate ellipsoid | $0 < e < 1$ | $-1 < k < 0$ | $f < R/2$ |
| Oblate ellipsoid | | $k > 0$ | |



Sag of aspheric surface

The aspheric sag equation is found simply by solving the conic equation for z .

Sag for a conic:

$$z = \frac{\frac{\rho^2}{R}}{1 + \sqrt{1 - (k + 1)\left(\frac{\rho}{R}\right)^2}}$$

Sag for even asphere:

$$z = \frac{\frac{\rho^2}{R}}{1 + \sqrt{1 - (k + 1)\left(\frac{\rho}{R}\right)^2}} + \alpha_4 \rho^4 + \alpha_6 \rho^6 + \dots$$

α 's = high-order
aspheric coefficients

R = vertex radius of curvature

Refs: Zemax manual (ver 13, 2013), p. 329

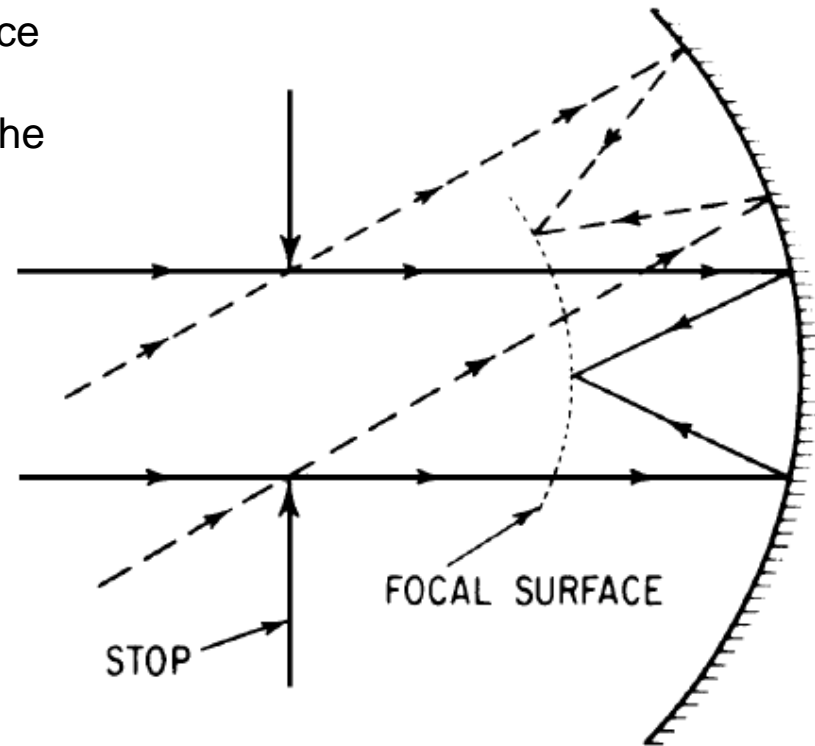
D. Malacara, Optical Shop Testing, 2nd ed., Appendix A

R. Shannon, Art and Science of Optical Design, sec. 7.1.6



Spherical Mirror with stop at center of curvature

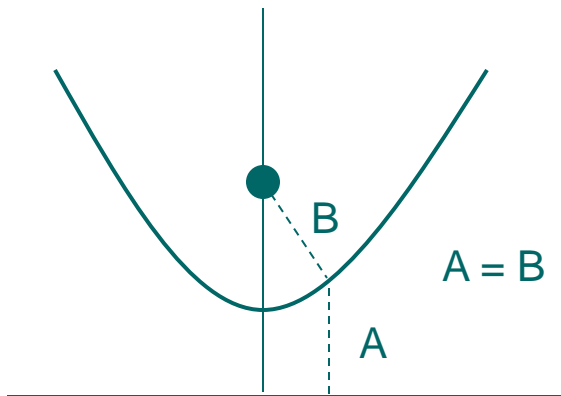
When the aperture stop is at the mirror's center of curvature, there are no “off-axis” rays and the image plane is a curved surface with radius equal to the mirror focal length. Coma and astigmatism are both zero, but the mirror still has spherical aberration.





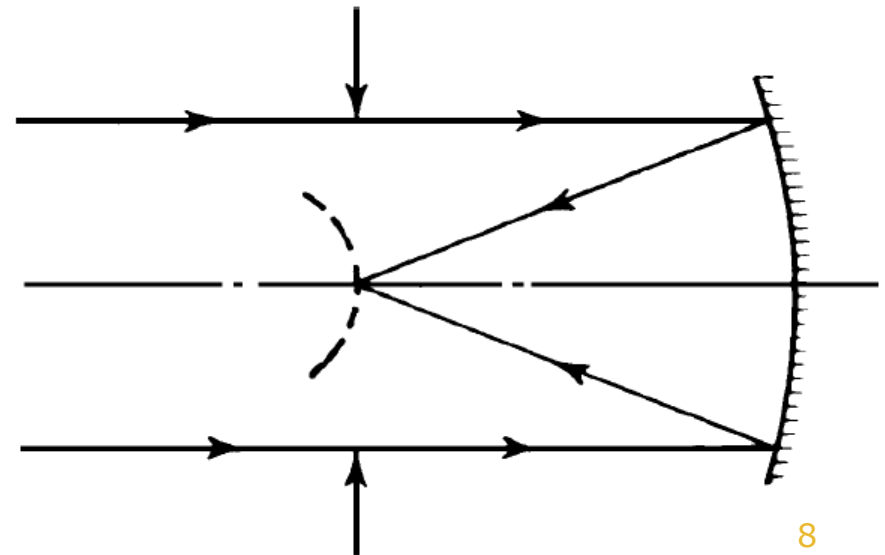
Paraboloidal Mirror

Parabola defined as the locus of points in a plane whose distance to the focus equals the distance to a fixed line (called the directrix).



Stigmatic imaging
for parallel light
(i.e., the “other”
focus is at infinity)

A parabolic mirror has zero astigmatism when the aperture stop is located at the focus. But, because spherical is also zero, stop shifts do not affect coma.

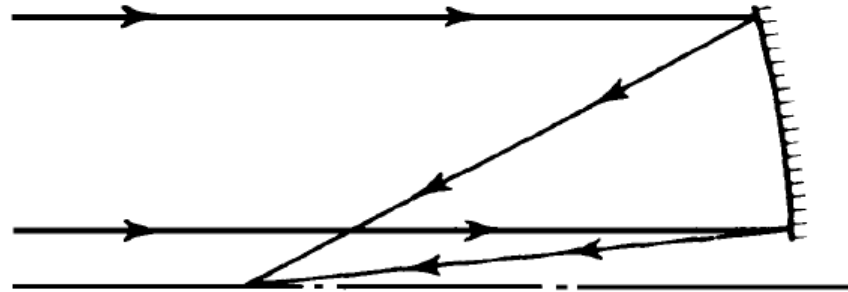




Paraboloidal one-mirror telescopes

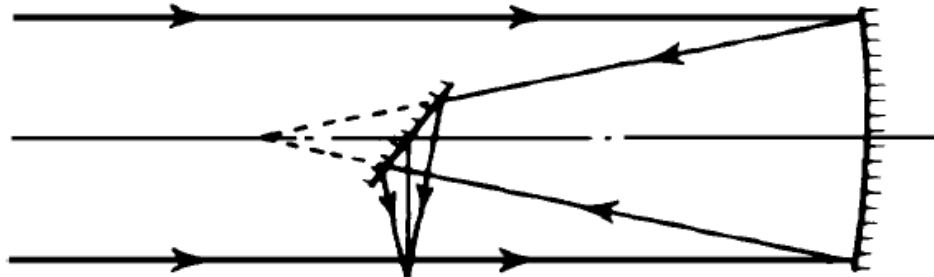
Herschellian

(off-axis parabola)



Newtonian

(high obscuration ratio)

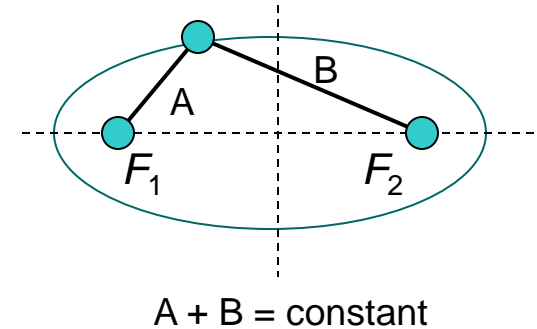


Field always limited by coma for parabolas.



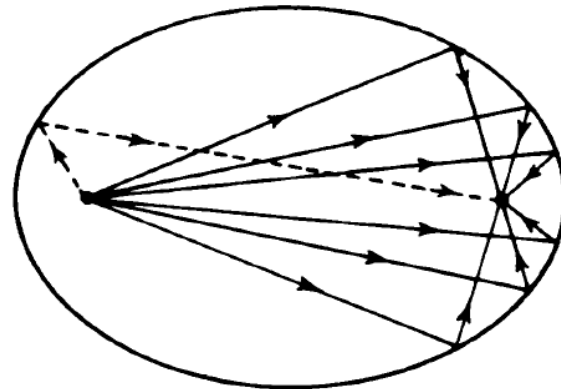
Ellipsoidal Mirror

Ellipse defined as the locus of points from which the sum of distances to two foci (F_1 and F_2) is a constant.



Optical property:

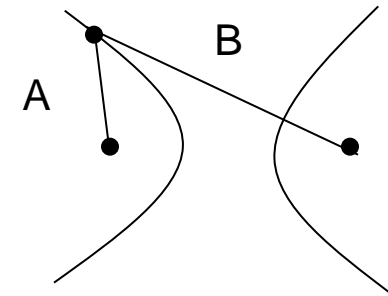
A ray passing through one focus is reflected to pass through the other focus.





Hyperboloid Mirror

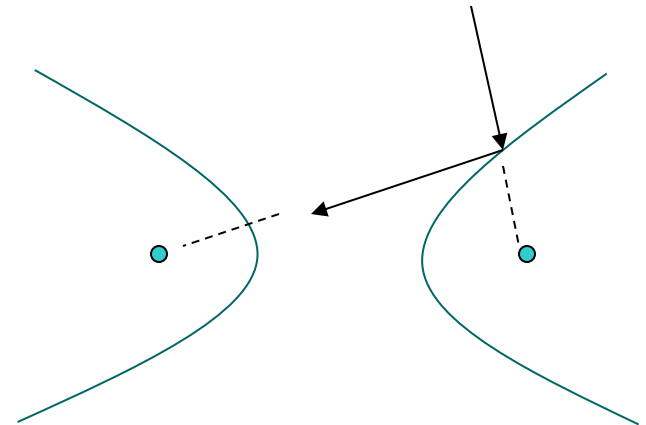
Hyperbola is defined as the locus of points for which the **|difference of distances|** to two foci is **constant** (points & foci in a common plane)



$$|A-B| = \text{constant}$$

Optical property:

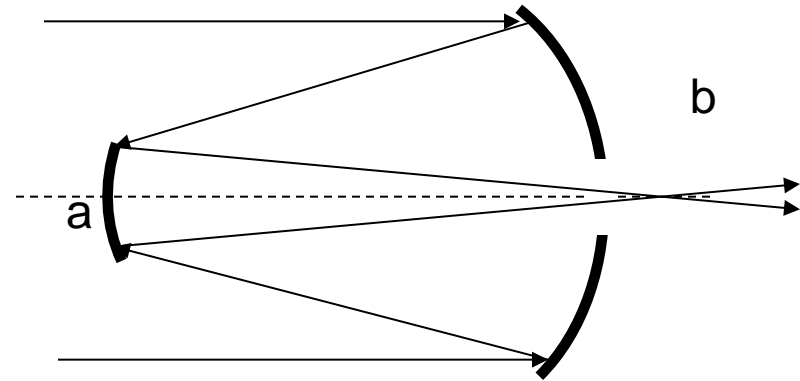
A ray that approaches a hyperbola from the side opposite a focus and pointing toward that focus reflects toward the other focus.





Two-mirror telescopes

$$\Phi = \Phi_a + \Phi_b - \frac{\Phi_a \Phi_b d}{n}$$



Relationship between primary-mirror and secondary-mirror conic constants (k), radius of curvature (R), marginal ray heights (y), and magnifications (m) for zero spherical aberration.

$$k_1 + 1 = \frac{y_2^4}{y_1^4} \frac{R_1^3}{R_2^3} \left[k_2 + \left(\frac{m-1}{m+1} \right)^2 \right]$$

D. Schroeder, *Astro. Optics*, p. 54 with m sign convention changed to match Smith's.

Ex: choose $k_2 = 0$ (sphere), solve for $k_1 = -0.7696$ (ellipse) ... *Dall-Kirkham*



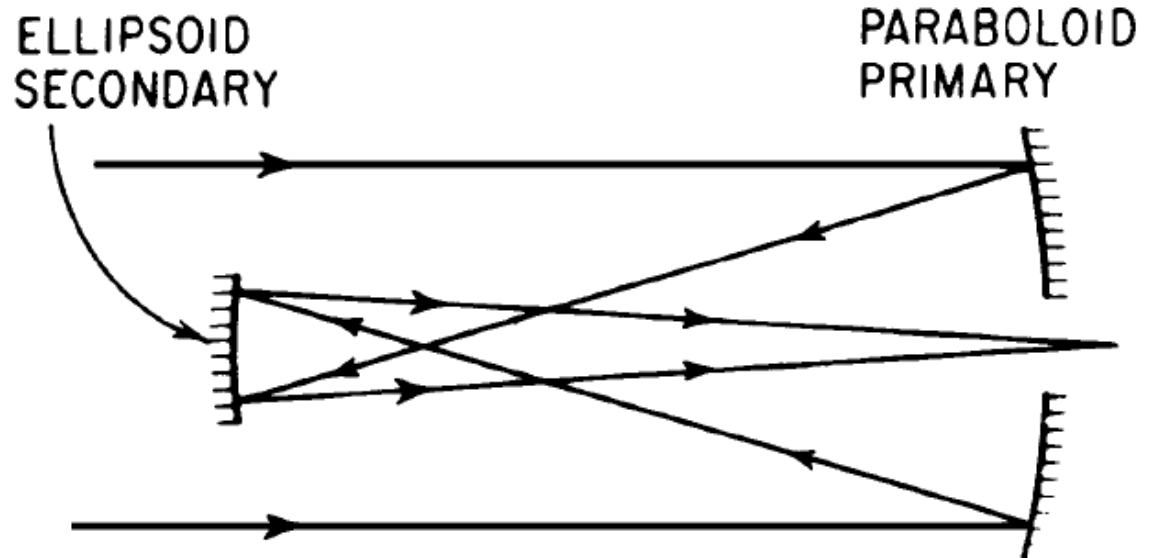
Gregorian

Primary mirror = parabola

Secondary = ellipsoid

Spherical = 0

Field limited by ...





Cassegrain

Primary mirror = parabola ($k = -1$)

Secondary = hyperbola

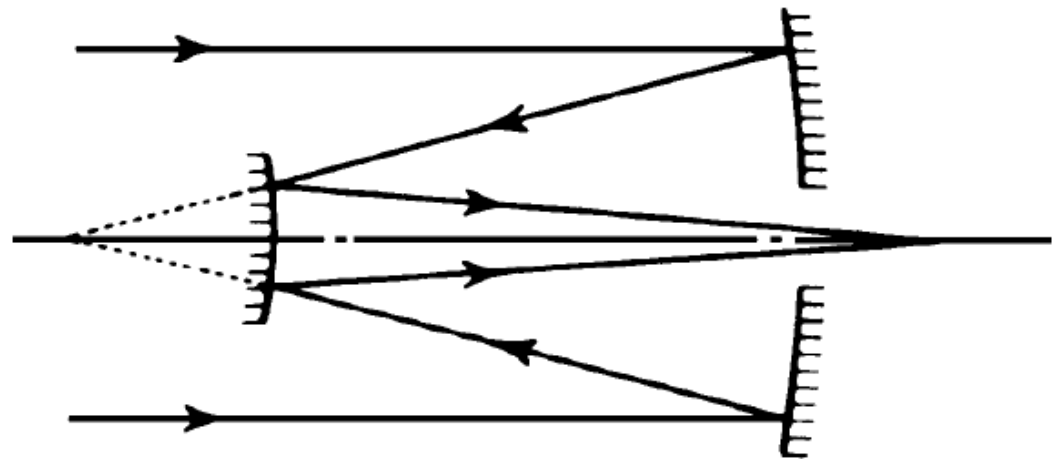
Spherical = 0

Astigmatism larger than for parabola alone,
but field usually limited by coma.

Shorter than optically
equivalent Gregorian

Cassegrain – France 1672

- + very compact
- convex secondary
difficult to test





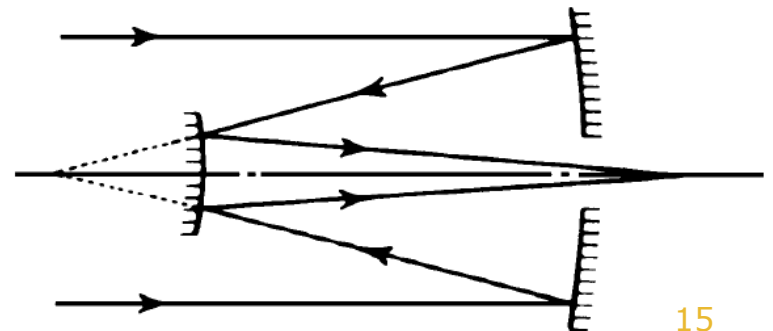
Dall-Kirkham

- Primary mirror = ellipse
- Secondary = sphere
- Spherical = 0 (on axis, obj at inf)
- Spherical secondary generates additional aberrations ...
- Field limited by coma, to about $\frac{1}{2}(m_2^2 + 1)$ times smaller than for classical cassegrain.

“poor man’s cassegrain”

(spherical 2ndary cheaper)

- + much easier to build and test than cassegrain (cheaper)
- + less sensitive alignment
- larger aberrations





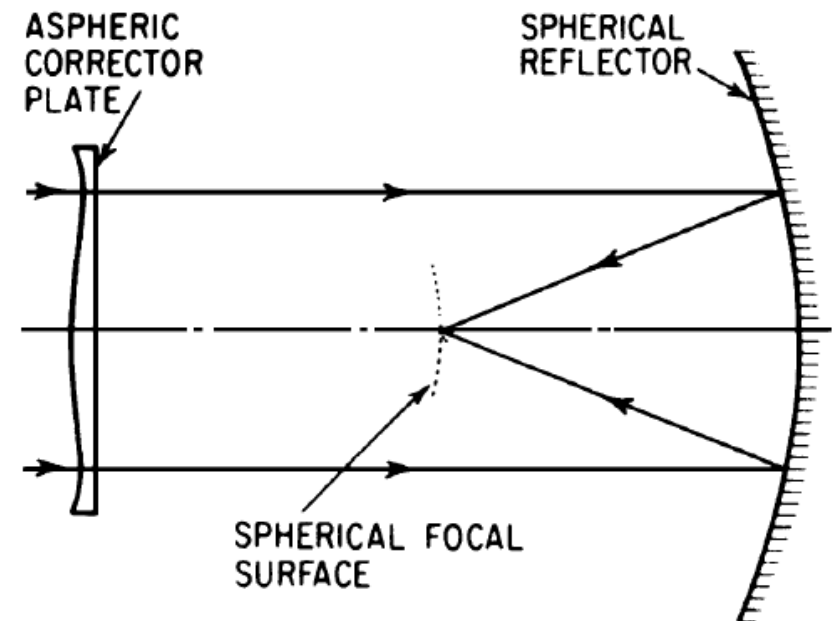
Schmidt

- “**Catadioptric**” system (combined reflector and refractor)
- Spherical mirror + aspheric refracting (corrector) plate at center of curvature
- Corrector plate at center leads to low off-axis aberrations
- Corrector plate minimizes spherical

+ very compact

+ corrector plate easier to build than parabolic primary

- images limited by spherochromatism and high-order astigmatism & spherical



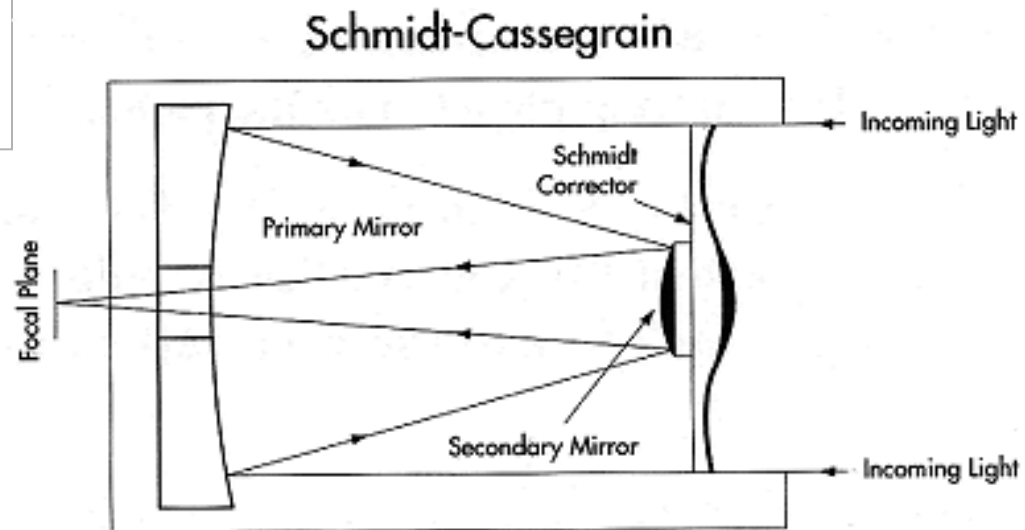


Schmidt-Cassegrain

Spherical primary and secondary plus aspheric corrector plate

- + very compact
- + easy to build and test
- images contain some coma, astigmatism, spherochromatism ...

Celestron C8 Photo





Maksutov

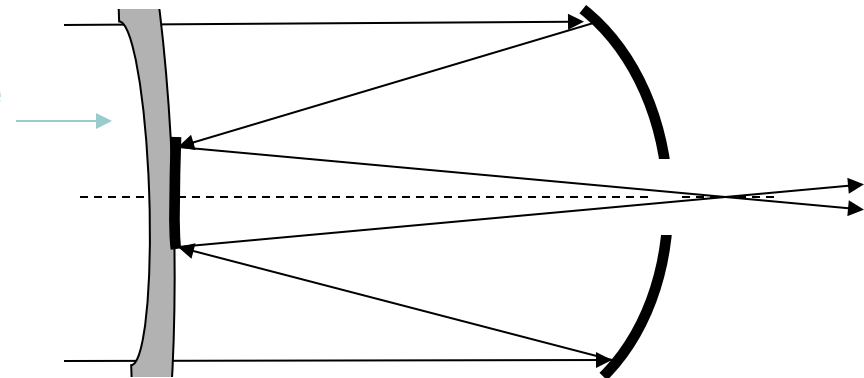
Spherical primary and secondary plus meniscus corrector plate
(sometimes elliptical primary ... ex. AstroPhysics)

- + very compact
- + easy to build and test
- images contain some coma, astigmatism, spherochromatism ...



Maksutov – Russia 1941

correcting plate
& secondary





Ritchey-Cretien

“Aplanatic Cassegrain”

~ Cassegrain with

primary mirror = hyperboloid

secondary = hyperboloid

(conic constant more negative than for classical cassegrain)

Spherical = coma = 0 (“aplanat”)

Slightly more astigmatism than cc, but zero coma allows larger FOV

Field limited by astigmatism (larger than coma-limit for cc because aberration symmetry)

+ wider field of view than cc

+ spherical = coma = 0

- expensive to build

- more difficult to test than other cassegrain types (Hubble = f/24 2.4 m R-C !)





Hubble Space Telescope

Ritchey-Cretien with 2.4-m primary and 0.3-m secondary (both hyperboloids)



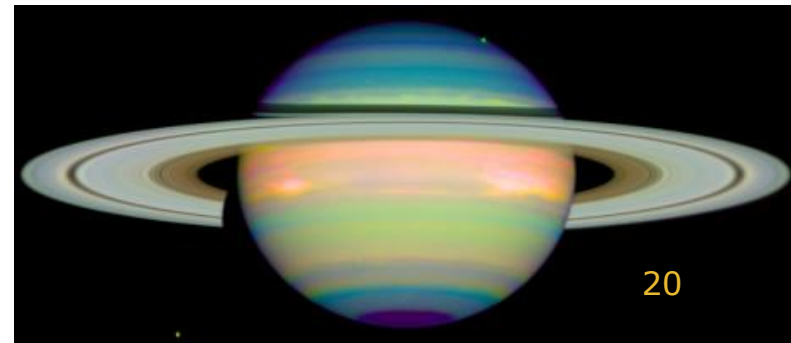
Before COSTAR



After

Hubble primary was ground to the wrong figure because of a slightly misplaced null lens during interferometric testing.

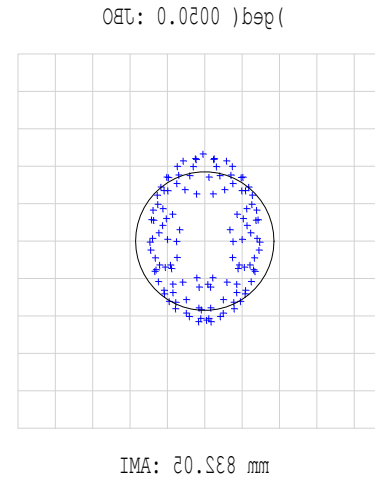
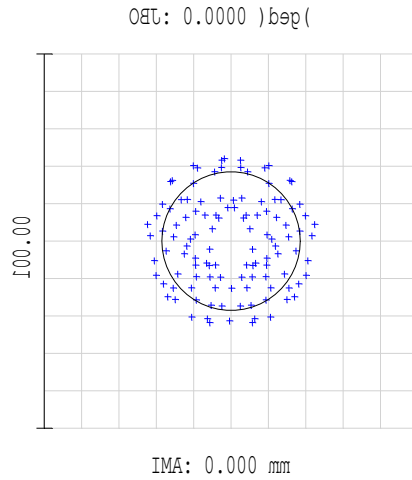
The result was large Spherical aberration.





+ 0.0000

Hubble spot diagrams – as designed



Source: IMA

Spot Diagram

Hubble Telescope
 3102\0\3
 Units are mm
 Field : 1
 RMS radius : 438.61
 RMS OEG : 228.22
 RMS radius : 100

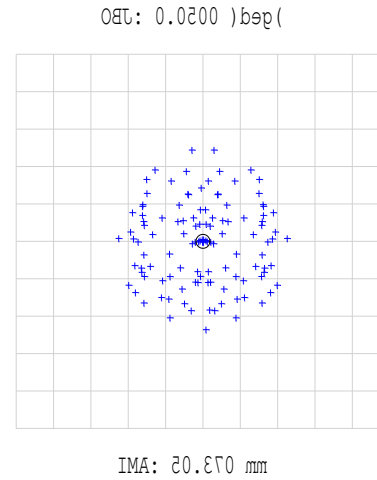
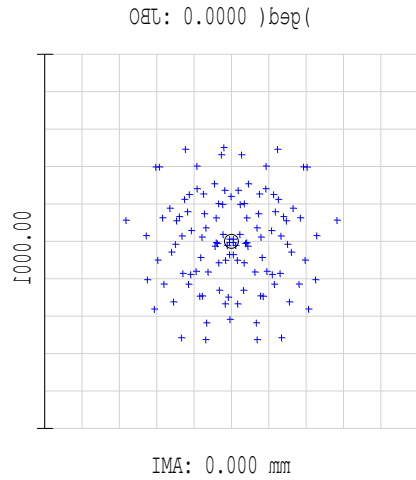
Airy radius: 18.25 μm
 2
 443.61
 222.32
 Var fehc : ecentered

HMZ(.nised)_1_elpbu
 I fo I noitarqifnoc



+ 0.8388

Hubble spot diagrams – as built



surface: IMA

spot diagrams

Hubble Telescope
 \3102\9\3 Units are um.
 Airy Radius: 18.81 um
 2 523.751
 1 608.961
 RMS radius : 223.092
 228.242
 Y-axis radius : 223.092
 X-axis radius : 228.242
 0001 : 1000

HMZ(.tlibsa)_l_elbbuH
 I fo l notarqifnoc

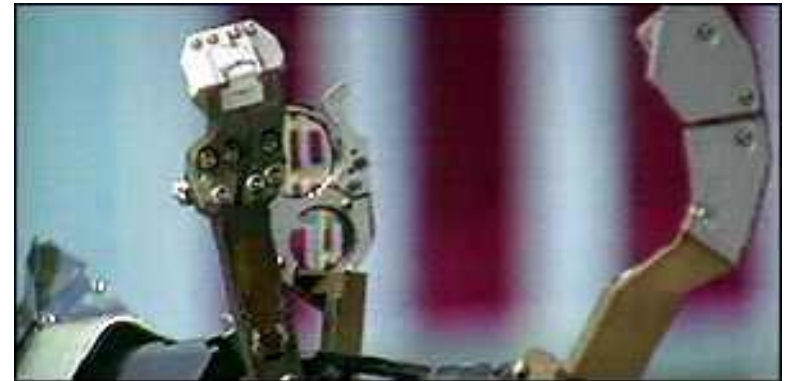


Hubble testing

The test interferometer was accidentally focused on a metal cover over the null corrector spacing rod instead of on the rod itself, as intended (this happened because a black paint chip fell off, revealing a metal glint point).



interferometric
null-lens test



COSTAR corrective
optics installed in 1993



See Jeff Hecht, "Saving Hubble," *Optics and Photonics News*, pp. 42-49, March 2013.

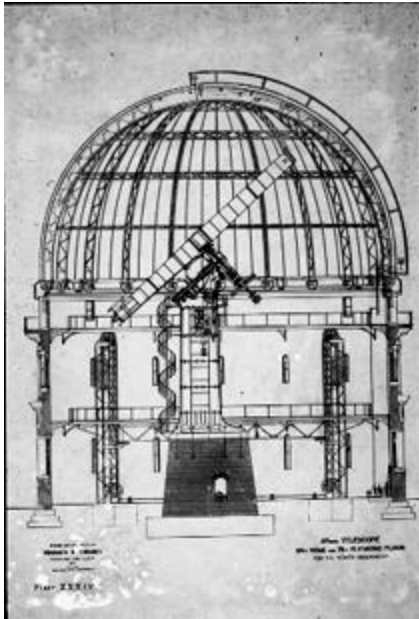
Images:
<http://news.bbc.co.uk/1/hi/sci/tech/638187.stm>



Large Telescopes

Modern mirrors can be made much larger than lenses, which would break or sag severely under their own weight.

- Largest refractor in the world is the 40-inch (1-m) telescope installed at Yerkes Observatory in Wisconsin by George Ellery Hale in 1897.



G. Hale 1928



<http://astro.uchicago.edu/tour/>



5-m Hale Telescope - Mt. Palomar





University of Arizona Mirror Lab

Pioneered technique of making huge mirrors on honeycomb substrates

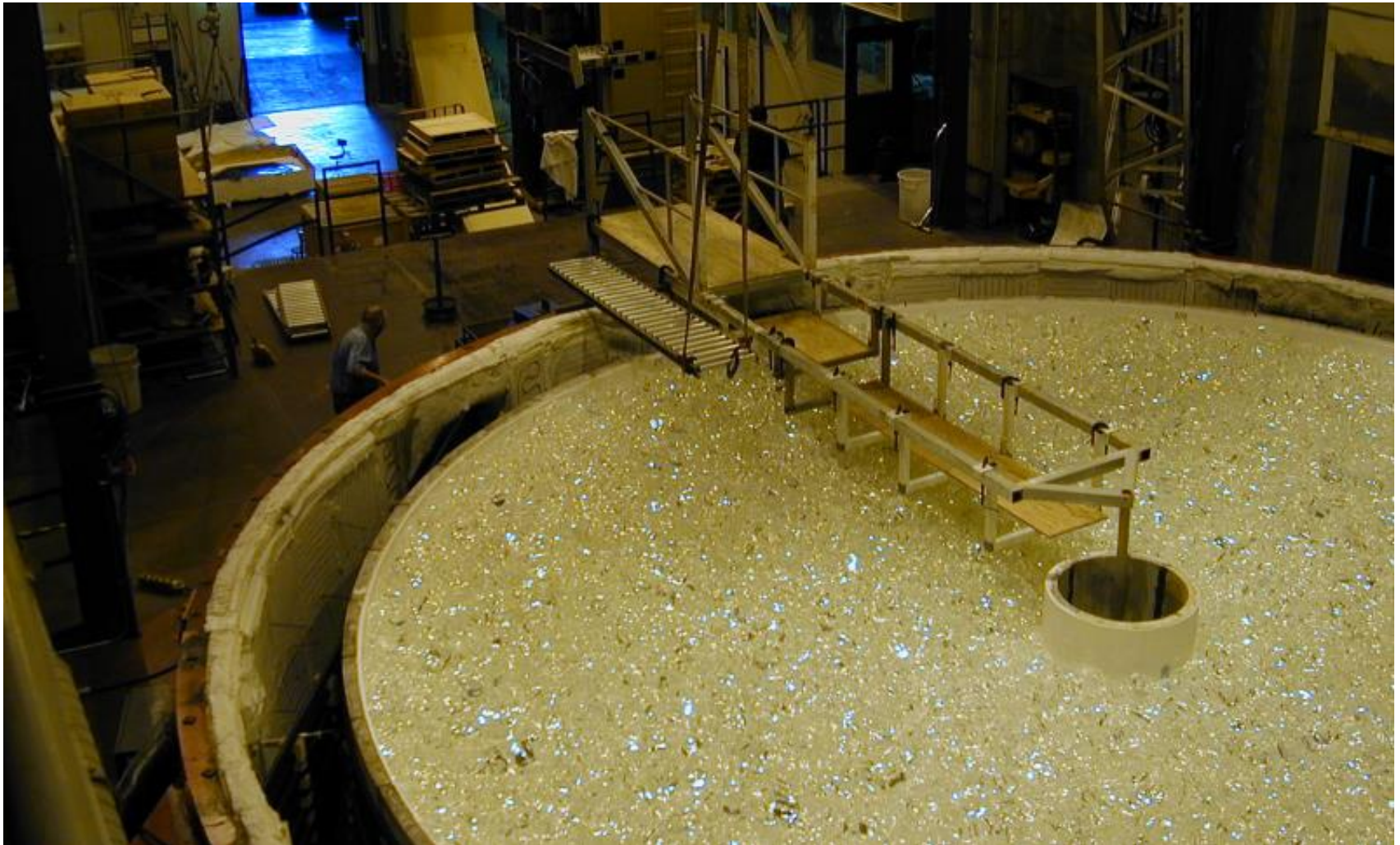
Largest casting to date = 8.4 m

- Borosilicate glass (used in making cooking dishes)
- honeycomb substrate provides strength with light weight
- First tried by Dr. Roger Angel in his backyard





University of Arizona 8.4-m mirror





University of Arizona 8.4-m mirror

Inserting hexagonal honeycomb structures for strength in a lightweight structure





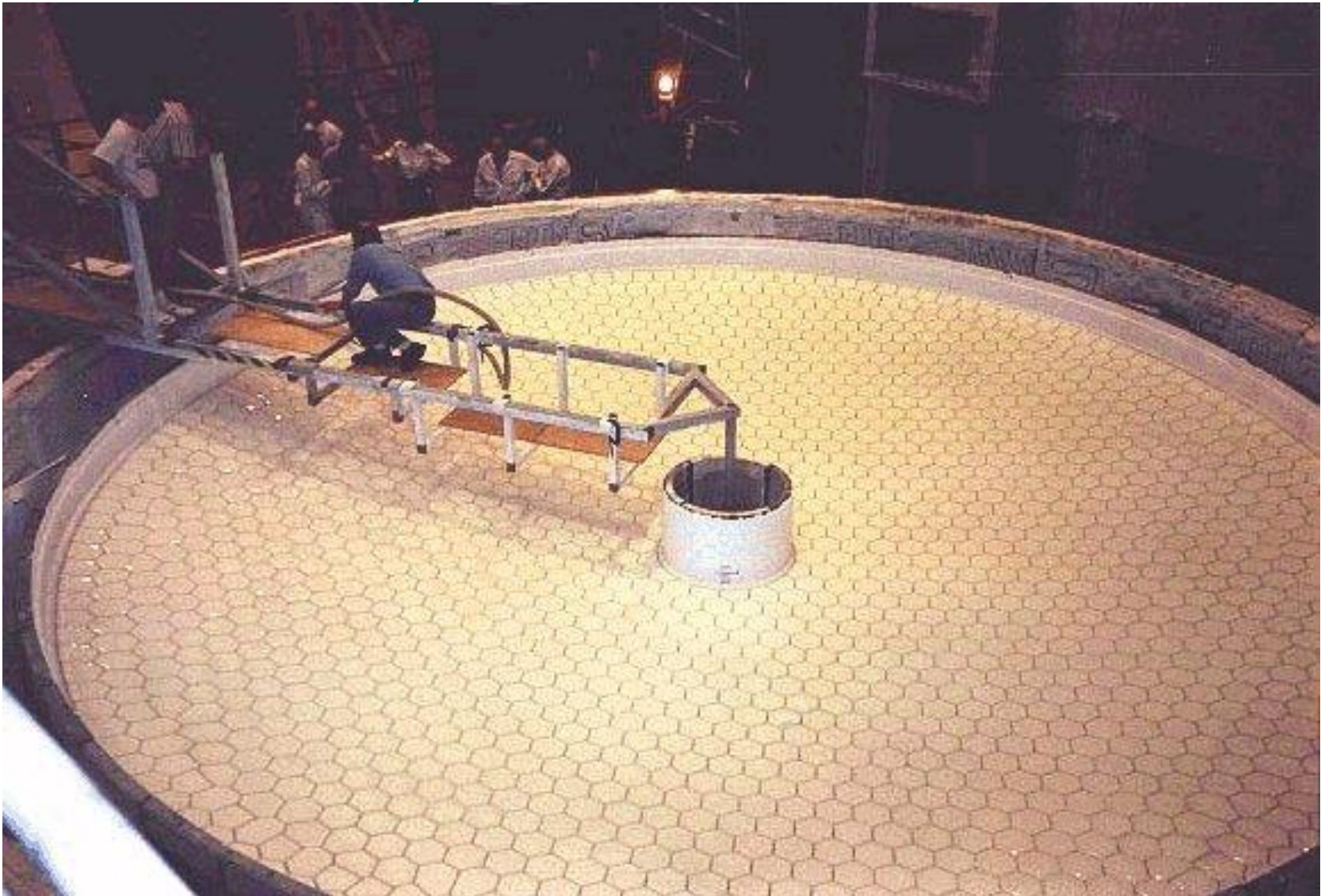
University of Arizona 8.4-m mirror



Grinding – part of a 6-year+ process to make the mirror



University of Arizona 8.4-m mirror





University of Arizona 8.4-m mirror !

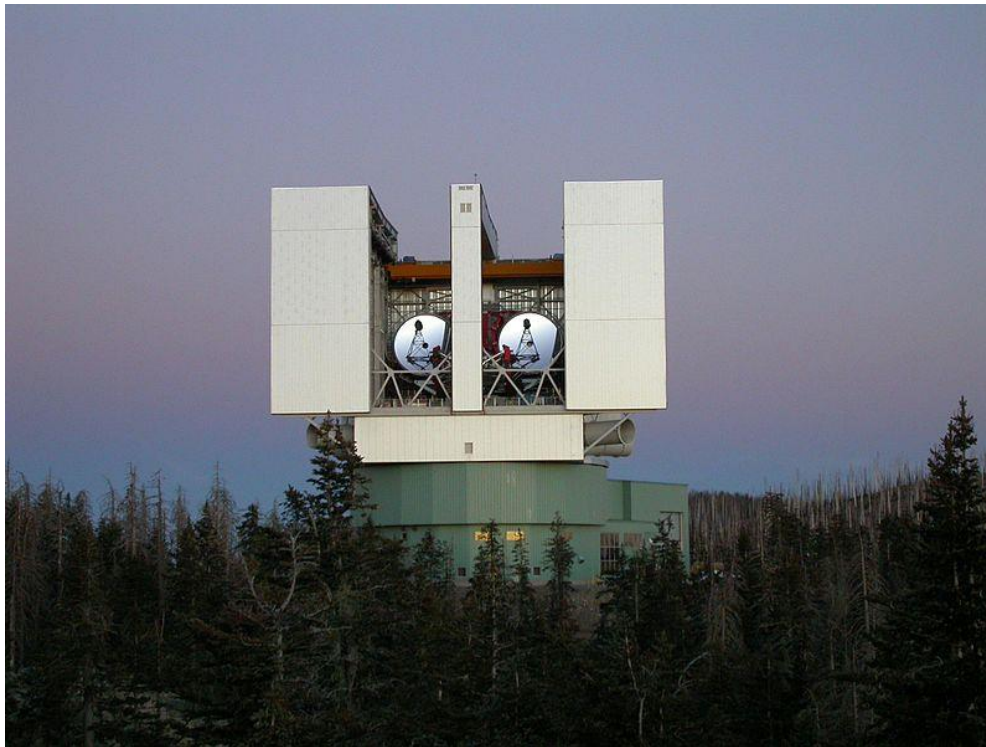
$D = 8.408 \text{ m}$ $F/1.142$



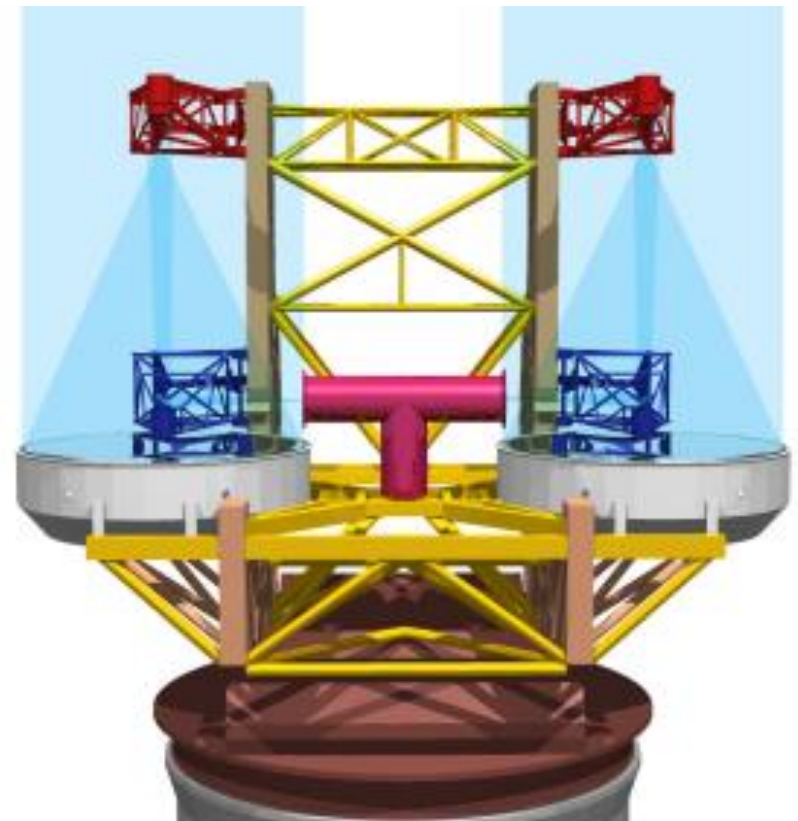


Large Binocular Telescope

Mt. Graham – southeast of Tucson, Arizona (first light 12 Oct. 2005)



Two $f/1.142$ mirrors:
8.4-m diameter, obscuration 0.889 m, center spacing 14.417 m
Collecting area equivalent to 11.8-m mirror
(diffraction-limited resolution equiv. to 22.8 m)



<http://mirrorlab.as.arizona.edu/index.php>

<http://medusa.as.arizona.edu/lbto/>



Giant Magellan Telescope

Las Campanas Observatory – NNE of La Serena, Chile in Atacama Desert (2018)



Three 8.4-m mirrors casted as of 2013

Aplanatic Gregorian

SEVEN off-axis 8.4-m primary mirrors
Collecting area equivalent to 21.4-m mirror
Diffraction-limited resolution equiv. to 24.5 m mirror

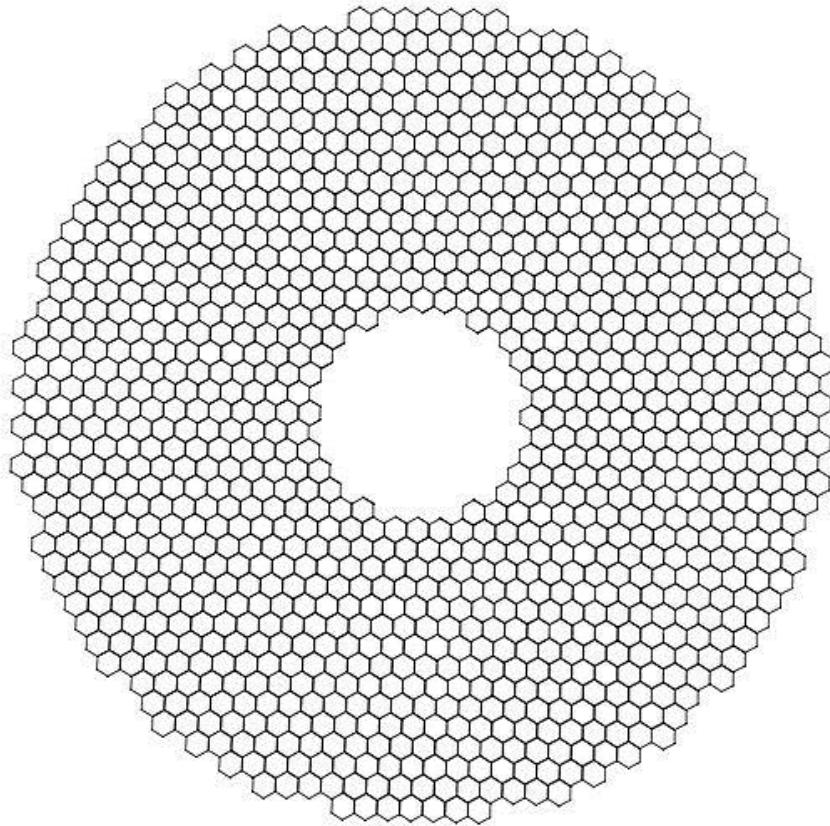
Seven off-axis adaptive secondary mirrors on hexapod mounts



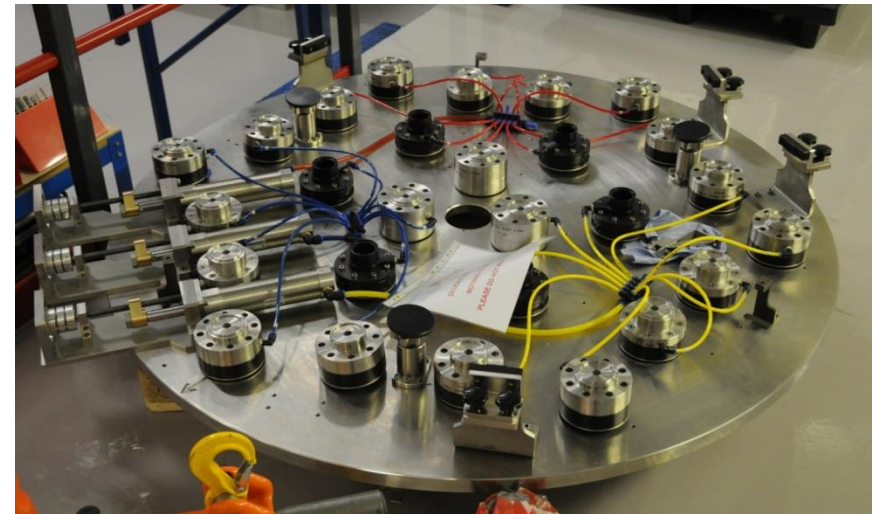


European Extremely Large Telescope

Cerro Armazones, Chile in Atacama Desert (2018)



~42-m primary mirror made of
~1.4-m adaptive segments



J. Shaw photo

<http://www.gmto.org/>



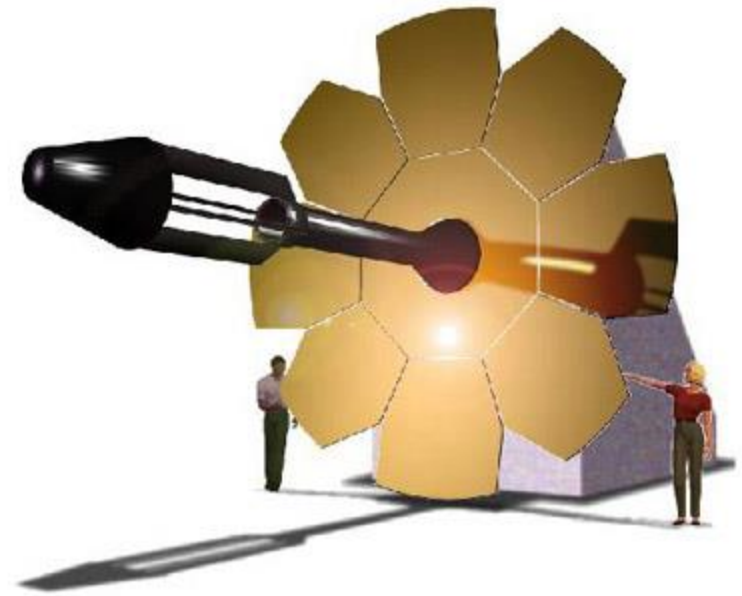
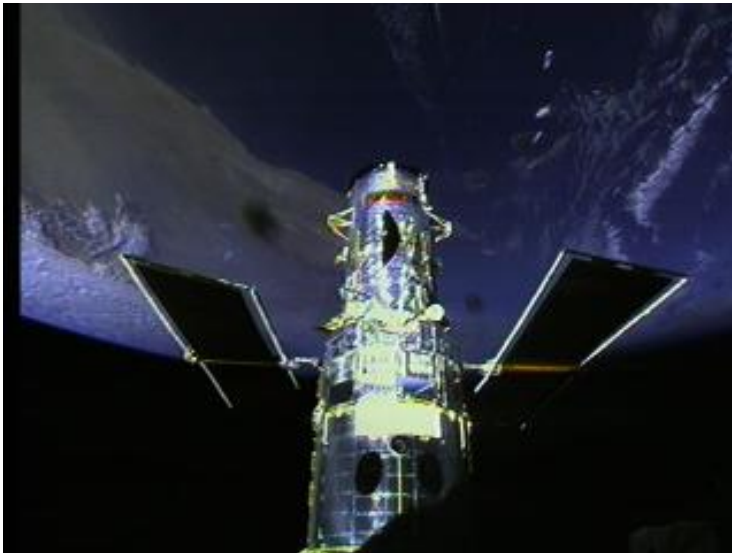
References

Daniel J. Schroeder, *Astronomical Optics*, Academic Press, 1987.

William L. Wolfe, *Introduction to Infrared System Design*, SPIE Press, 1996.

Warren J. Smith, *Modern Optical Engineering*, 3rd ed., McGraw-Hill, 2000.

W. Welford, *Aberrations of the Symmetrical Optical System*, Adam-Hilger, 1986.





Appendix

Additional information for further study



Spherical Mirror

- Surface = section of circle ($k = 0$)
- Simplest reflector
- highest local radius of curvature
- Perfect imaging at 1:1 conjugates ($m = 1$)
- Undercorrected spherical and off-axis aberrations at non-unity conjugates
- For distant objects, spherical aberration is $\sim 1/8$ that of an equivalent lens bent at ‘minimum bending’

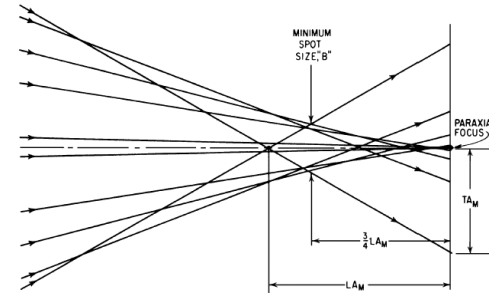
Tx 3rd-order spherical for spherical mirror with magnification m , at paraxial focus (Smith p. 475 & Schroeder p. 47)

$$TSA3 = \frac{(m-1)^2}{m+1} \frac{\rho^3}{2R^2}$$



Spherical aberration for a spherical mirror

For 3rd-order spherical, the minimum geometrical spot size is at $\frac{3}{4}$ of the distance from the paraxial focus to the marginal focus.



At this location, the minimum blur spot diameter is equal to $\frac{1}{2}$ the minimum transverse spherical:

$$B = \frac{(m-1)^2}{(m+1)} \frac{\rho^3}{4R^2}$$

Divide by $f = R/2$ to obtain the angular blur in radians:

$$\beta = \frac{(m-1)^2}{(m+1)} \frac{\rho^3}{2R^3}$$

Use $F\# = f/(2\rho)$, $f = R/2$, and $m = 0$ to rewrite this as
(need to add high-order aberrations below about $f/2$)

$$\beta = \frac{1}{128(F\#)^3} = \frac{NA^3}{16}$$



Spherical mirror off-axis aberrations (stop not at center)

For distant object, the 3rd-order coma and astigmatism contributions are found by using $n' = -n$ in the surface-aberration equations from Smith's chapter 10 (which are just scaled versions of the Seidel sums) ...

Sagittal 3rd-order coma

$$CC^* = \frac{\rho^2 (R - \bar{l}) \bar{u}}{2R^2} = \frac{(R - \bar{l}) \bar{u}}{32(F\#)^2}$$

Longitudinal 3rd-order astigmatism
(1/2 separation of S and T fields)

$$AC^* = \frac{(\bar{l} - R)^2 \bar{u}^2}{4R}$$

Longitudinal Petzval

$$PC = \frac{\bar{u}^2 R}{4} = \frac{h_i^2}{2f}$$

Note: when the Stop is at distance R , $CC^* = AC^* = 0$.

In these equations ρ is the normalized aperture radial coordinate, R is the mirror radius of curvature, h_i is the image height, f is the focal length, \bar{l} is the distance from vertex to aperture stop, and \bar{u} is the chief-ray angle at the image.



Spherical mirror with stop at the mirror

When the aperture stop is located at the mirror, the minimum angular blur sizes are:

3rd-order sagittal coma $\beta_c = \frac{-\bar{u}}{16(F\#)^2}$ radians

3rd-order astigmatism at best focus $\beta_a = \frac{\bar{u}^2}{2(F\#)}$ radians



Spherical mirror with stop in between

When the aperture stop is located somewhere between the mirror and the center of curvature, we can simply add the minimum blur spot sizes from The equations for spherical blur, coma blur, and astigmatism blur, adjusted for the following scaling:

1. coma = astigmatism = 0 when stop is at center of curvature
2. coma varies linearly with distance of stop from center of curvature
3. astigmatism varies quadratically with distance of stop from center...

Keep in mind that the appropriately scaled β sum is still only approximate (uses only 3rd-order aberrations) and is angular blur in units of radians.